

# Multi Grid

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# Outline

What's in Wednesday's Directory?

References

MG for Integral Equations

MG for Elliptic PDEs

Exercises

- ▶ This lecture
- ▶ Revised transport lecture
- ▶ Two papers
- ▶ MATLAB examples

Tuesday's directory has also been revised.

# References I

- ▶ W. L. BRIGGS, V. E. HENSON, AND S. MCCORMICK, A Multigrid Tutorial, 2nd edition, Society for Industrial and Applied Mathematics, Philadelphia, 2000.
- ▶ A. BRANDT, Multilevel adaptive solutions to boundary value problems, *Math. Comp.*, 31 (1977), pp. 333–390.
- ▶ W. HACKBUSCH, Multi-Grid Methods and Applications, vol. 4 of Springer Series in Computational Mathematics, Springer-Verlag, New York, 1985.
- ▶ S. MCCORMICK, Multilevel Adaptive Methods for Partial Differential Equations, SIAM, Philadelphia, 1989.

## References II

- ▶ K. E. ATKINSON, Iterative variants of the Nyström method for the numerical solution of integral equations, Numer. Math., 22 (1973), pp. 17–31.
- ▶ C. T. KELLEY, A fast multilevel algorithm for integral equations, SIAM J. Numer. Anal., 32 (1995), pp. 501–513.
- ▶ C. T. KELLEY, Multilevel source iteration accelerators for the linear transport equation in slab geometry, Trans. Th. Stat. Phys., 24 (1995), pp. 679–708.

# MG for Integral Equations

Fredholm 2nd kind integral equation on  $C[0, 1]$ ,

$$u(x) = (Ku)(x) + g(x) = \int_0^1 k(x, y)u(y) dy + g(x).$$

Here

- ▶  $k$  and  $g$  are given continuous functions.
- ▶  $u \in C[0, 1]$  is the unknown.

We assume that  $I - K$  is nonsingular.

## Atkinson (73) - Brakhage (60) Method

Notation:

- ▶ Sequence of quadrature rules: nodes  $\{x_j^m\}_{j=1}^{N_m}$  and weights  $\{w_j^m\}_{j=1}^{N_m}$ .
- ▶ Operators:  $K_m(u)(x) = \sum_{j=1}^{N_m} k(x, x_j^m)u(x_j^m)w_j^m$
- ▶ Note that  $K_m$  is defined on  $C[0, 1]$ .

The sequence  $\{K_m\}$  is collectively compact and converges strongly to  $K$ .

Defining all operators on  $C[0, 1]$  is easier than thinking of sequences of problems and operators.

## Solving $u - K_m u = f$

First solve the finite-dimensional system for the function values at the nodes

$$u(x_i^m) - \sum_{j=1}^{N_m} k(x_i^m, x_j^m) u(x_j^m) w_j^m = f(x_i^m)$$

(with GMRES, for example).

Then recover  $u$  with the Nyström interpolation

$$u(x) = f(x) + \sum_{j=1}^{N_m} k(x, x_j^m) u(x_j^m) w_j^m.$$

## Some Facts from Functional Analysis

Given  $\rho > 0$  there is  $l_0$  such that if  $L \geq l \geq l_0$  the operator

$$B_l^L = I + (I - K_l)^{-1} K_L$$

satisfies

$$\|I - B_l^L(I - K_L)\| \leq \rho.$$

NOTE!!  $(I - K_l)^{-1}$  won't work! We will fix that by changing  $K_l$ .

## Atkinson-Brakhage iteration

Solve  $u - K_L u = g$  on a fine mesh by using  $B_I^L$  as a preconditioner to fixed point iteration.

$$u_+ = u_c - B_I^L(u_c - K_L u_c - g).$$

Cost: Two fine-mesh calls to  $K_L$ .

$r_c = g - (I - K_L)u_c$  and  $B_I^L r_c = r_c + (I - K_I)^{-1} K_L r_c$ .  
Evaluate  $(I - K_I)^{-1} w$  via GMRES.

# Implementation

You have seen how to compute  $K_L$ . The last part is the action of  $B_I^L$  on a function  $u$ .

- ▶ Evaluate  $K_L u$ .
- ▶ Solve  $w - K_I w = K_L u$  as above.

## A More Efficient Way

- ▶ Suppose you have a compact operator  $K$  and
- ▶ can compute operator-function products.
- ▶ Suppose you have a projection  $P_l$  onto a finite dimensional subspace  $V_l$  and
- ▶  $P_l$  converges strongly to  $I$ .
- ▶ Then  $KP_L$  converges to  $K$  in the operator norm and
- ▶  $(I - KP_L)^{-1} \rightarrow (I - K)^{-1}$  in norm.

## Faster Two-Grid Method

$$u_+ = u_c - (I - KP_I)^{-1}((I - K)u_c - g)$$

How to evaluate  $(I - KP_I)^{-1}w$ .

- ▶ Solve the finite dimensional problem

$$w_I - P_I K P_I w_I = P_I g$$

- ▶ Let

$$w = g - K P_I w_I.$$

Nested Iteration:  $K_i = KP_i; K = K_L$ 

Algorithm `gridnest`( $g, u, \{K_i\}, l, L$ )

Solve  $u^l - \mathcal{K}_l u^l = g$  on the coarse level.

Set  $u = u^l$ .

**for**  $m = l + 1, \dots, L$  **do**

$$u = u - (I - KP_l)^{-1}((I - K)u - g)$$

**end for**

## Remarks

- ▶ Analysis shows that  $\|K_I - K_L\|$  is small for sufficiently large  $I$  and all  $L > I$  (including  $L = \infty$ , the continuous problem).
- ▶ The algorithm can be realized by using piecewise linear interpolation, evaluation at grid points, and **full weighting** for the coarse-to-fine transfer.

## Easy example

$$u(x) - \frac{1}{2} \int_0^1 \sin(x-y)u(y) dy = \cos(x)$$

I have already discretized this into the form  $(I - K)u = f$ , so the hard job is the intergrid transfers.

- ▶ Coarse to fine: simple interpolation.
- ▶ Fine to coarse: repeated applications of full weighting.
- ▶ And here's some MATLAB ...

# Nested Iteration

## Opimal approach

- ▶ Start on coarse grid  $l$ ; solve to high accuracy.
- ▶ Interpolate to fine grid. Solve with multilevel method.
- ▶ **Keep the coarse level at  $l$  for the entire solve!** $k$
- ▶ You'll need one iteration/level to maintain accuracy to truncation error if the coarse mesh is fine enough.

## Example: Source Iteration in Transport Theory

- ▶ Let  $\phi - K_L\phi = S$  be the fully discrete problem at level  $L$ .
- ▶ Map from level  $L$  to level  $l < L$  by full weighting.
- ▶ That map has the same properties as a continuous projection.
- ▶ Map from  $l$  to  $L$  by interpolation.
- ▶ So  $K_l = I_l^L K_L I_l^L$  will do the job.

# Multigrid

We showed that a Jacobi iteration for the central difference discretization  $A^h u^h = f^h$  of

$$-u'' = f, 0 < x < 1$$

with  $u(0) = u(1) = 0$

- ▶ damped the error in the high frequencies fast, and
- ▶ the error in the low frequencies very slowly,
- ▶ so the iteration is a **smoother**,

which explains the methods poor convergence properties.

MG exploits this by attacking the smooth components on coarser grids.

# Discrete Laplacian I

Recall that

$$A^h = \frac{1}{h^2} \begin{pmatrix} 2 & -1 & 0 & \dots & 0 & 0 \\ -1 & 2 & -1 & ,0 & \dots & 0 \\ 0 & -1 & 2 & -1 & , \dots & 0 \\ \vdots & \ddots & \ddots & \ddots & \ddots & \vdots \\ 0 & \dots, & ,0, & -1 & 2 & -1 \\ 0 & \dots, & \dots,, & 0 & -1 & 2 \end{pmatrix}$$

and

$$f^h = (f(x_1), \dots, f(x_N))^T.$$

## Discrete Laplacian II

The eigenvectors of  $A^h$  are  $\{u_n\}_{n=1}^N$  where

$$u_n = (\xi_1^n, \dots, \xi_N^n)^T \text{ and } \xi_i^n = \sqrt{2/h} \sin(ni\pi h),$$

with eigenvalues

$$\lambda_n = h^{-2} 2(1 - \cos(\pi nh)) = \pi^2 n^2 + O(h^2) \text{ for small } n$$

We proved this with elementary trigonometry.

## A better smoother

Consider the **damped Jacobi** iteration:

$$x_{n+1} = (1 - \omega)x_n - \omega D^{-1}(L + U)x_n + D^{-1}b$$

with iteration matrix  $M_{DJ}^h = (1 - \omega)I + \omega M_{JAC}$ , where

$$M_{JAC}^h = -D^{-1}(L + U).$$

So  $\omega = 1$  is Jacobi and  $\omega = 0$  is nothing.

What's the optimal  $\omega$  if we want to damp the high-frequency terms?

## Eigen-analysis

We showed that the eigenvalues of  $M_{JAC}$  were

$$\mu_n = 1 - (h^2/2)\lambda_n$$

with the same eigenvectors as  $A^h$ . Similarly

$$\begin{aligned}\mu_n &= (1 - \omega) + \omega(1 - (h^2/2)\lambda_n) = 1 - (h^2/2)\omega\lambda_n \\ &1 - \omega(1 - \cos(\pi nh))\end{aligned}$$

## Optimal $\omega$

Now suppose  $n \geq N/2$  (high frequency), then

$$\mu_n = 1 - \omega(1 - \cos(\pi nh)) = 1 - \omega + \omega \cos(\pi nh)$$

so

$$1 - 2\omega \leq \mu_n \leq 1 - \omega.$$

Minimize  $|\mu_n|$  to see that the optimal value of  $\omega$  is  $2/3$ . So

$$|\mu_n| \leq 1/3 \text{ for } N/2 \leq n \leq N - 1$$

# Idealized Two-Grid Method: I

Notation:

- ▶  $h = 1/(N - 1)$ ,  $N - 1$  even.
- ▶  $\Omega^h$ : space of grid functions with step size  $h$   
so  $u^h \in \Omega^h$
- ▶  $I_h^{h'}$ : intergrid transfer from  $\Omega^h \rightarrow \Omega^{h'}$
- ▶  $P_H^h$  and  $P_L^h$  project onto low and high frequencies.

## Idealized Two-Grid Method: II

**For Now:** We will define  $I_h^{2h}$  and  $I_{2h}^h$  by Fourier truncation.

$$u = \sum_{n=1}^{N-1} \alpha_n u_n^h \in \Omega^h$$

$$I_h^{2h} u = \sum_{n=1}^{(N-1)/2} \alpha_n u_n^{2h} \in \Omega^{2h} = I_h^{2h} P_L^h u$$

and

$$w = \sum_{n=1}^{(N-1)/2} \alpha_n u_n^{2h} \in \Omega^{2h}$$

$$I_{2h}^h w = \sum_{n=1}^{(N-1)/2} \alpha_n u_n^h \in \Omega^h$$

Note:  $I_h^{2h} P_L A^u = A^{2h} I_h^{2h} P_L u$  and  $I_{2h}^h I_h^{2h} = P_L$

## Idealized Two-Grid Method: III

Simple method:

- ▶ Let  $u_c \in \Omega^h$
- ▶ Take one weighted Jacobi iteration on  $\Omega^h$  to obtain  $u_{1/2}$
- ▶ Compute the residual  $r^h = f^h - A^h u_{1/2}$
- ▶ Compute the **coarse grid correction**  $w = (A^{2h})^{-1} I_h^{2h} r^h$
- ▶  $u_+ = u_{1/2} + I_{2h}^h w$

## Idealized Two-Grid Method: IV

If we did everything on  $\Omega^h$ , then

$$w = A^{-1}r = A^{-1}(f^h - A^h u) = A^{-1}f^h - u = u^h - u$$

so

$$u + w = u^h$$

solves the problem. We will show that the two grid method does pretty well.

## Idealized Two-Grid Method: V

Apply weighted Jacobi and let  $e = u^h - u$ .

$$\|P_H e_{1/2}\| \leq (1/3)\|P_H e_0\|$$

so

$$\begin{aligned} r^h &= f^h - A^h u_{1/2} = P_H(f^h - A^h u_{1/2}) + P_L(f^h - A^h u_{1/2}) \\ &= A^h P_H e_{1/2} + P_L A^h e_{1/2}. \end{aligned}$$

We are just about done because ...

## Idealized Two-Grid Method: VI

The coarse grid correction eliminates the low frequency errors.

- ▶  $I_h^{2h} P_L A^h e_{1/2} = A^{2h} P_L e_{1/2}$ , so
- ▶  $(A^{2h})^{-1} I_h^{2h} P_L A^h e_{1/2} = I_h^{2h} P_L e_{1/2}$ ,
- ▶ and  $w = I_{2h}^h I_h^{2h} P_L e_{1/2} = P_L e_{1/2}$ .

Hence

$$e_+ = e_{1/2} - P_L e_{1/2} = P_H e_{1/2}$$

and hence  $\|e_+\| \leq (1/3)\|e_c\|$

## Algorithmic Description

`two_grid`( $h, u_c^h, u_+^h, A, f$ )

Smooth once to obtain  $u_{1/2}^h$  from  $u_c^h$

$$r^h = f^h - A^h u_{1/2}^h$$

Solve  $A^{2h} w = I_h^{2h} r^h$  exactly.

$$u_+^h = u_{1/2}^h + I_{2h}^h w$$

## Algorithmic Description: MG

Multigrid replaces the coarse mesh solve with two-grid iteration.

$h$ : desired (finest) grid;  $H$ : coarsest grid.

`multi_grid( $h, H, u_c^h, u_+^h, A, f$ )`

**if**  $h = H$  **then**

Solve  $A^h u^h = f^h$  exactly.

**else**

Smooth once to obtain  $u_{1/2}^h$  from  $u_c^h$

$$r^h = f^h - A^h u_{1/2}^h$$

`multi_grid( $2h, H, 0, w, A, I_h^{2h} r^h$ )`

$$u_+^h = u_{1/2}^h + I_{2h}^h w$$

**end if**

This is a V-cycle. The convergence rate remains unchanged.

## Observations

- ▶ MG is faster than two-grid because the exact solves live only on the coarsest mesh.
- ▶ This method only works for the simplest problem where
  - ▶ we can connect the eigenfunctions of  $A$  to those of  $M$
  - ▶ the intergrid transfers have the same eigenfunctions
  - ▶  $I_{2h}^h I_h^{2h} = P_L^h$
- ▶ In general, we have none of those things.

## Realistic Intergrid Transfers

- ▶ Nested grids:  $x_i^h = ih$ ;  $x_{2i}^h = x_i^{2h}$
- ▶  $I_{2h}^h$ : linear interpolation;

$$(I_{2h}^h v^{2h})_{2i} = v_i^{2h} \text{ and } (I_{2h}^h v^{2h})_{2i-1} = (v_i^{2h} + v_{i-1}^{2h})/2.$$

- ▶  $I_h^{2h}$ : full weighting

$$(I_h^{2h} v^h)_i = (v_{2i-1}^h + 2v_{2i}^h + v_{2i+1}^h)/4.$$

except on the boundary.

# Matrix Representation: I

- ▶  $N = 2^p - 1$  internal grid points
- ▶ Internal grid points:  $\{x_i^h\}_{i=1}^N$ ;  $h = 1/(N + 1)$
- ▶ Boundary grid points:  $x_0^h = 0$ ,  $x_{N+1}^h = 1$
- ▶  $I_{2h}^h$  is  $1 + (N - 1)/2 \times N$
- ▶ Other option: restriction by injection:  $(I_h^{2h}v)_i = v_{2i}$

## Matrix Representation: II

$$I_{2h}^h = \frac{1}{2} \begin{pmatrix} 1 & 0 & \dots & 0 & 0 & 0 \\ 2 & 0 & \dots & 0 & 0 & 0 \\ 1 & 1 & \dots & 0 & 0 & 0 \\ 0 & 2 & \dots & 0 & 0 & 0 \\ 0 & 1 & 1 & \dots & 0 & 0 \\ 0 & 0 & 2 & \dots & 0 & 0 \\ 0 & 0 & 1 & 1 & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & \dots & 1 & 1 \\ 0 & 0 & 0 & \dots & 0 & 2 \\ 0 & 0 & 0 & \dots & 0 & 1 \end{pmatrix}, I_h^{2h} = \frac{1}{2} (I_{2h}^h)^T.$$

## Comments

- ▶ Interpolatory intergrid transfers have aliasing effects.
- ▶ You have to smooth an extra time.
- ▶ But you get the convergence rate  $1/9$  instead of  $1/3$  you'd expect.

## Intergrid Transfers for 1-D: $I_{2h}^h$

```
w = I_{2h}^h(u)
N = h^{-1} - 1; h_2 = 2h; N_2 = h_2^{-1} - 1
for i = 0 : N_2 do
    w_{2i} = u_i
    w_{2i+1} = (u_i + u_{i+1})/2
end for
```

## Intergrid Transfers for 1-D: $I_h^{2h}$

```
w = I_h^{2h}(u)
N = h^{-1} - 1; h_2 = 2h; N_2 = h_2^{-1} - 1
w_0 = u_0; w_{N_2+1} = u_{N+1}
for i = 1 : N_2 do
    w_i = (u_{2i-1} + 2u_{2i} + u_{2i+1})/4
end for
```

# Smoother

Notation: Smooth  $\nu$  times at level  $h$  with initial iterate  $u^h$   
 $u^h \leftarrow S(u^h, f^h, \nu)$

```
for  $is = 1 : \nu$  do  
  for  $i = 1 : N$  do  
     $w_i = h^2(f_i^h + u_{i-1}^h + u_{i+1}^h)/2$   
  end for  
  for  $i = 1 : N$  do  
     $u_i^h = (1 - \omega)u_i^h + \omega w_i$   
  end for  
end for
```

## The V-cycle

$$u^h \leftarrow V^h(u^h, f^h)$$

**if**  $h = H$  **then**

Solve  $A^h u^h = f^h$  to high accuracy.

**else**

$$u^h \leftarrow S(u^j, f^h, \nu_1)$$

Compute  $f^{2h} = I_h^{2h}(f^h - A^h u^h)$

$$u^{2h} = 0$$

$$u^{2h} \leftarrow V^{2h}(u^{2h}, f^{2h})$$

$$u^h \leftarrow u^h + I_{2h}^h u^{2h}$$

$$u^h \leftarrow S(u^h, f^h, \nu_2)$$

**end if**

## Comments

- ▶ Typical choice  $\nu_1 = \nu_2 = 1$ .
- ▶ V-cycle for high-order term is a good preconditioner for PDEs.
- ▶ V-cycle reduces the error by fixed amount independent of  $h$ .
- ▶ We consider very easy problems and simple grids
- ▶ Solve step for  $\Omega^H$  can be iterative.
- ▶ There's more, much more
  - ▶ Other smoothers (Gauss-Seidel, ...)
  - ▶ Algebraic multigrid (AMG)
  - ▶ MG for Navier-Stokes, combustion ...
  - ▶ Non isotropic flows/grids

## Matlab Example: Coarse Mesh Solve

You can avoid the exact solve on the coarse mesh with no loss.  
This example shows that

- ▶ Convergence rate is independent of  $h$
- ▶ You can simply do a few (24) smoothing steps at the coarse level
- ▶ You are not restricted to Poisson's equation

# Helmholtz in 1-D

$$u'' + \sigma u = f; u(0) = u(1) = 1.$$

Arrange things so the exact solution is

$$u(x) = e^x \sin(\pi x).$$

## Comparisons from Coarse\_Solve\_Test directory

- ▶ Several fine levels:  $h = 2^{-l}$ ,  $5 \leq l \leq 10$
- ▶ Exact solve vs 24 smooths on Coarse mesh.

## Nested Iteration or Full Multigrid (FMG): $h = 2^{-p}H$

Solve  $A^H u^H = f^H$  on  $\Omega^H$ ;  $h = H$ .

**for**  $il = 1 : p$  **do**

$$h = h/2$$

$$u^h = I_{2h}^h u^{2h}$$

$$u^h \leftarrow V^h(u^h, f^h)$$

**end for**

# Data Structures

You do not want to repeatedly allocate storage for the solutions on the various grids. You can avoid this by creating a large structure which preallocates the storage.  
Some of the examples in the **FMG\_Test** do this.

## Comments

- ▶ Optimal complexity (see next slide)
- ▶ Complexity leads to debugging tool
- ▶ FMG is a **solver**. If error is  $O(h^2)$  then one V-cycle per level will suffice.
- ▶ A V-cycle can be a **preconditioner**

## Complexity Example: I

- ▶  $A^h$  discrete Laplacian in  $R$ ,  $N$  interior grid points
- ▶ Cost of smoother, mat-vec, intergrid transfer =  $3N$
- ▶ Cost of residual, correction =  $N$
- ▶ Cost of solve at level  $H$  is  $V_0$
- ▶ Then, cost of V-cycle at level  $h_j = 2^{-j}H$  is

$$V_j = 3(N_j(\nu_1 + \nu_2) + 2N_j) + 2N_j + V_{j-1}$$

where  $N_j = (1/h_j)^2 - 1 \approx 2N_{j-1}$ .

## Complexity Example: II

So, if  $\nu_1 = \nu_2 = 1$ , then

$$\begin{aligned} V_j &= 14N_j + V_{j-1} = 14 \sum_{l=1}^j N_l + V_0 \\ &\leq 28N_j + V_0 \end{aligned}$$

Cost of FMG for  $h = 2^{-p}H$ ;  $N = N_p$

$$\begin{aligned} FMG_i &\leq \sum_{l=0}^p N_l + V_l \leq \sum_{l=0}^p 29N_l + pV_0 \\ &\leq 58N + pV_0 = O(N). \end{aligned}$$

## Testing Complexity

In one dimension, you should see the computing time double if  $h \rightarrow h/2$ .

In practice, MATLAB-sized problems take so little time that you will find this hard to measure.

## Exercises

- ▶ Solve the transport equation with mutlgrid in space using the modified Atkinson-Brakhage method.
- ▶ Write a MG V-cycle for the 2-D Poisson Equation. Test it as both a solver and a preconditioner within the `k1pde2ddemo.m`.
- ▶ How does the performance of your V-cycle preconditioner for the convection-diffusion problem in `k1pde2ddemo.m` change if you use two V-cycles instead of one?